

Panel Discussion: Probing New Physics with Double Beta Decays beyond the Tonne Scale

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Probing New Physics with $\beta\beta$

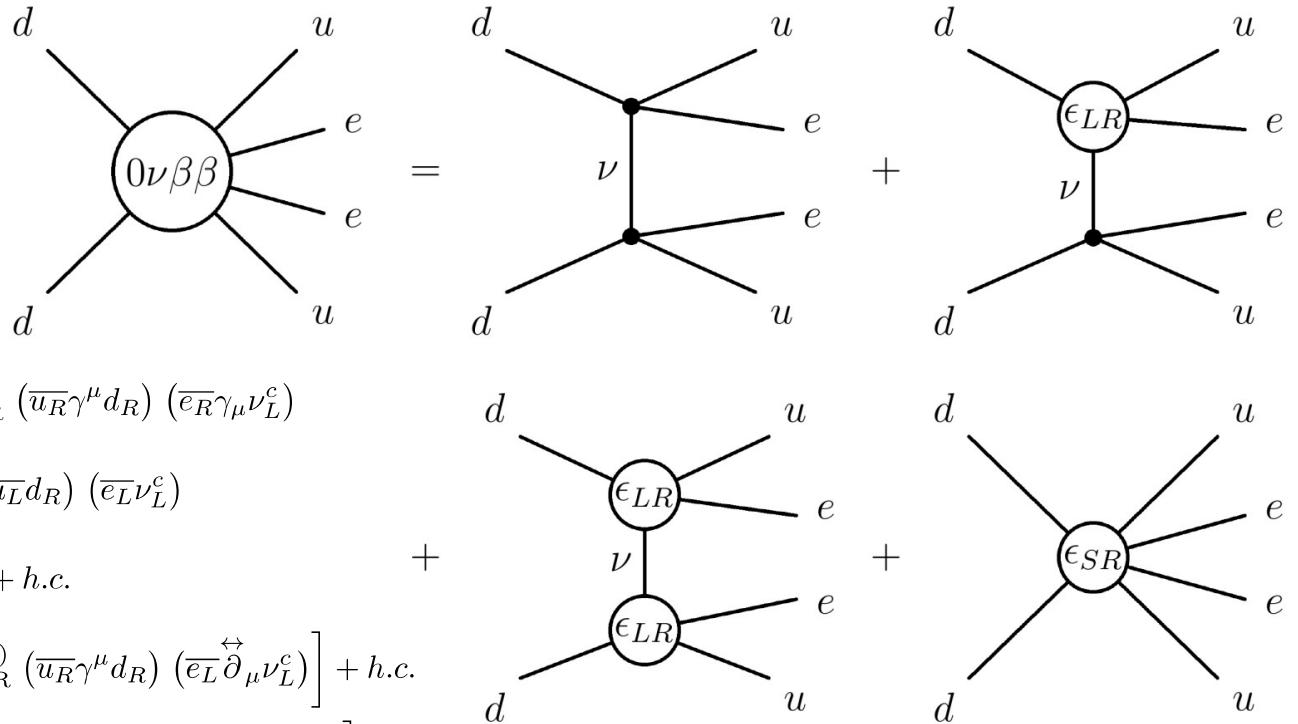


- Majorana masses \leftrightarrow LNV \leftrightarrow neutrinoless double beta decay ($0\nu\beta\beta$)
- UV physics – which mechanism?
 - LG, M. Lindner, O. Scholer: *Distinguishing among $0\nu\beta\beta$ Mechanisms*, in preparation
- Interplay of different contributions – non-trivial, complex formulae
 - J. de Vries, LG, O. Scholer: *$0\nu\beta\beta$ Automation Tool*, in preparation
- searches for $0\nu\beta\beta$ decay \rightarrow large amount of $2\nu\beta\beta$ background decay data collected $\sim 10^5 - 10^6$ events
- New Physics in $2\nu\beta\beta$ decay?
 - right-handed currents \rightarrow F. F. Deppisch, LG, F. Simkovic: [2003.11836](#), [PRL 125](#)
 - ν self-interactions \rightarrow F. F. Deppisch, LG, W. Rodejohann, X. Xu: [2004.11919](#), [PRD 102](#)
 - sterile neutrinos \rightarrow P. D. Bolton, F. F. Deppisch, LG, F. Simkovic: [2011.13387](#)

Non-Standard $0\nu\beta\beta$

- effectively, a variety of different mechanisms beyond the standard scenario may contribute to $0\nu\beta\beta$ (1208.0727, 1708.09390, 1806.02780, 1806.06058, 2009.10119)

- more descriptions:
epsilon basis
vs. Wilson
coefficients



$$\mathcal{L}_{\Delta L=2}^{(6)} = \frac{2G_F}{\sqrt{2}} \left[C_{VL}^{(6)} (\overline{u}_L \gamma^\mu d_L) (\overline{e}_R \gamma_\mu \nu_L^c) + C_{VR}^{(6)} (\overline{u}_R \gamma^\mu d_R) (\overline{e}_R \gamma_\mu \nu_L^c) \right. \\ \left. + C_{SL}^{(6)} (\overline{u}_R d_L) (\overline{e}_L \nu_L^c) + C_{SR}^{(6)} (\overline{u}_L d_R) (\overline{e}_L \nu_L^c) \right. \\ \left. + C_T^{(6)} (\overline{u}_L \sigma^{\mu\nu} d_R) (\overline{e}_L \sigma_{\mu\nu} \nu_L^c) \right] + h.c.$$

$$\mathcal{L}_{\Delta L=2}^{(7)} = \frac{2G_F}{\sqrt{2}v} \left[C_{VL}^{(7)} (\overline{u}_L \gamma^\mu d_L) (\overline{e}_L \overleftrightarrow{\partial}_\mu \nu_L^c) + C_{VR}^{(7)} (\overline{u}_R \gamma^\mu d_R) (\overline{e}_L \overleftrightarrow{\partial}_\mu \nu_L^c) \right] + h.c.$$

$$\mathcal{L}_{\Delta L=2}^{(9)} = \frac{1}{v^5} \sum_i \left[\left(C_{i,R}^{(9)} (\overline{e}_R e_R^c) + C_{i,L}^{(9)} (\overline{e}_L e_L^c) \right) \mathcal{O}_i + C_i^{(9)} (\overline{e} \gamma_\mu \gamma_5 e^c) \mathcal{O}_i^\mu \right]$$

Distinguishing $0\nu\beta\beta$ Mechanisms

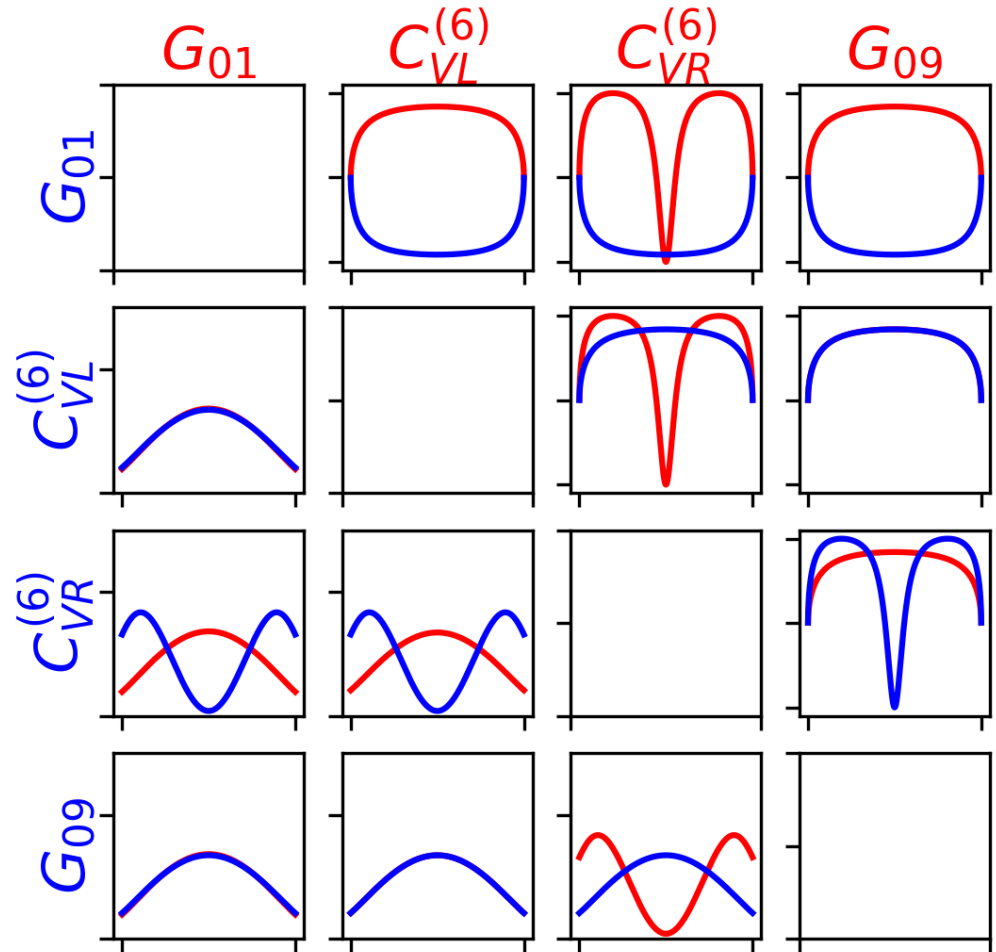
- phase-space observables – electron energy spectra, angular correlation
- comparison with other $\beta\beta$ modes $\rightarrow \beta+\beta+$, $EC\beta+$, $ECEC$ - typically suppressed
- decay rate ratios for different isotopes $R^{\mathcal{O}_i}(^AX) \equiv \frac{T_{1/2}^{\mathcal{O}_i}(^AX)}{T_{1/2}^{\mathcal{O}_i}(^{76}\text{Ge})} = \frac{|\mathcal{M}^{\mathcal{O}_i}(^{76}\text{Ge})|^2 G^{\mathcal{O}_i}(^{76}\text{Ge})}{|\mathcal{M}^{\mathcal{O}_i}(^AX)|^2 G^{\mathcal{O}_i}(^AX)}$
 - \rightarrow ratio of half-lives = ratio of NMEs x ratio of PSFs, the unknown coupling drops out
 - distinguishing 2 specific operators quantified using $R_{ij}(^AX) = \frac{R^{\mathcal{O}_i}(^AX)}{R^{\mathcal{O}_j}(^AX)}$
- applied to the “master formula” framework of 1806.02780 - V. Ciriigliano, W. Dekens, J. de Vries, M.L. Graesser, E. Mereghetti: *A neutrinoless double beta decay master formula from effective field theory*, JHEP 12
 - PSFs \rightarrow 4 distinguishable groups of operators
 - ratios: in principle 12 distinguishable groups of operators
 - main issue: unknown low energy constants (LECs) \rightarrow large uncertainties \rightarrow LQCD computations vital

Distinguishing $0\nu\beta\beta$ Mechanisms - PSFs

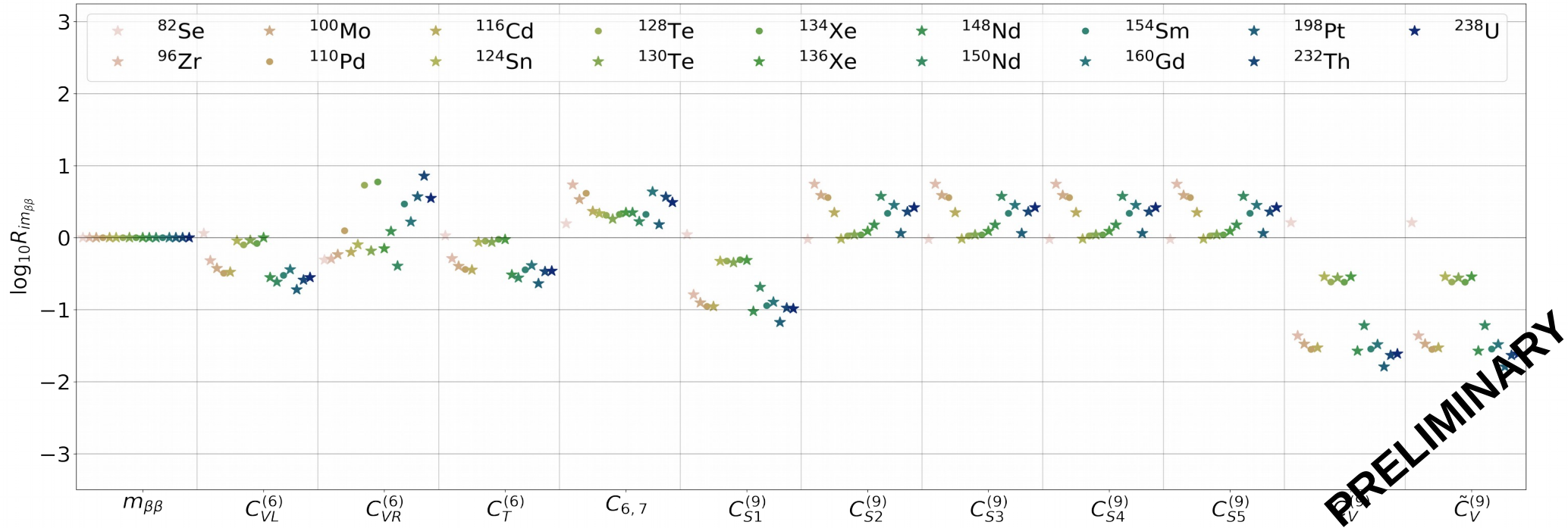
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$$\mathcal{L}_{\Delta L=2}^{(7)} = \frac{2G_F}{\sqrt{2}v} \left[C_{VL}^{(7)} (\bar{u}_L \gamma^\mu d_L) (\bar{e}_L \overset{\leftrightarrow}{\partial}_\mu \nu_L^c) \right. \\ \left. + C_{VR}^{(7)} (\bar{u}_R \gamma^\mu d_R) (\bar{e}_L \overset{\leftrightarrow}{\partial}_\mu \nu_L^c) \right] + h.c.$$

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Distinguishing $0\nu\beta\beta$ Mechanisms - Ratios



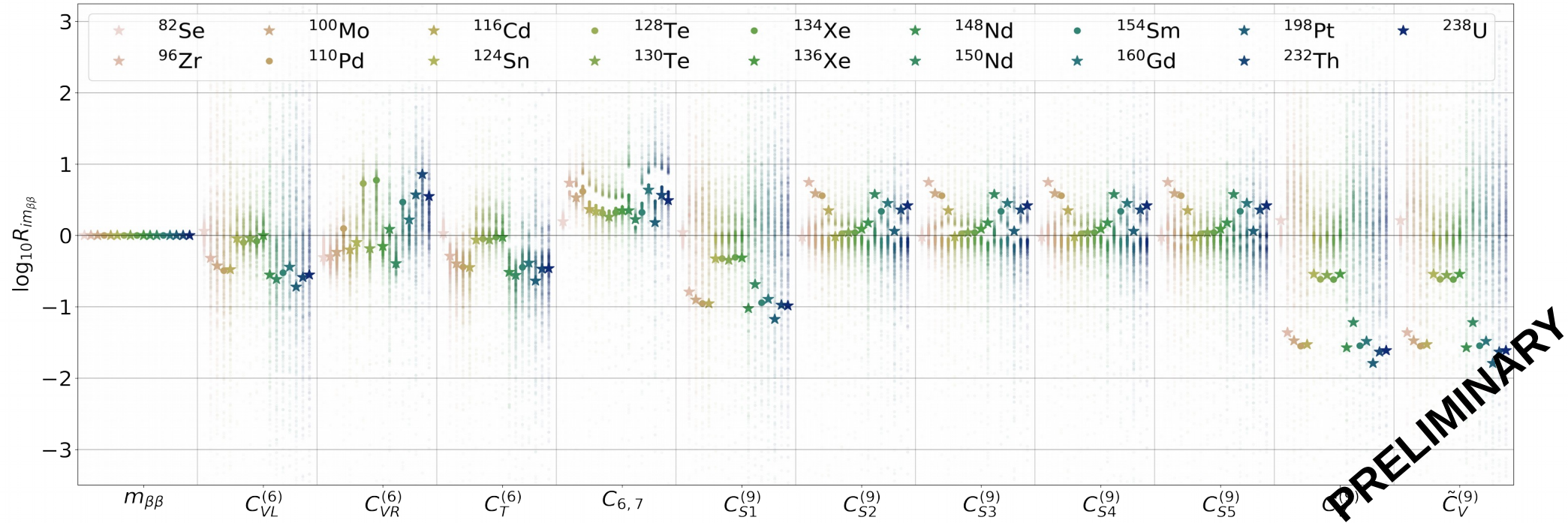
PRELIMINARY

$$\mathcal{L}_{\Delta L=2}^{(6)} = \frac{2G_F}{\sqrt{2}} \left[C_{VL}^{(6)} (\overline{u_L} \gamma^\mu d_L) (\overline{e_R} \gamma_\mu \nu_L^c) + C_{VR}^{(6)} (\overline{u_R} \gamma^\mu d_R) (\overline{e_R} \gamma_\mu \nu_L^c) \right. \\ \left. + C_{SL}^{(6)} (\overline{u_R} d_L) (\overline{e_L} \nu_L^c) + C_{SR}^{(6)} (\overline{u_L} d_R) (\overline{e_L} \nu_L^c) \right. \\ \left. + C_T^{(6)} (\overline{u_L} \sigma^{\mu\nu} d_R) (\overline{e_L} \sigma_{\mu\nu} \nu_L^c) \right] + h.c.$$

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Distinguishing $0\nu\beta\beta$ Mechanisms - Ratios



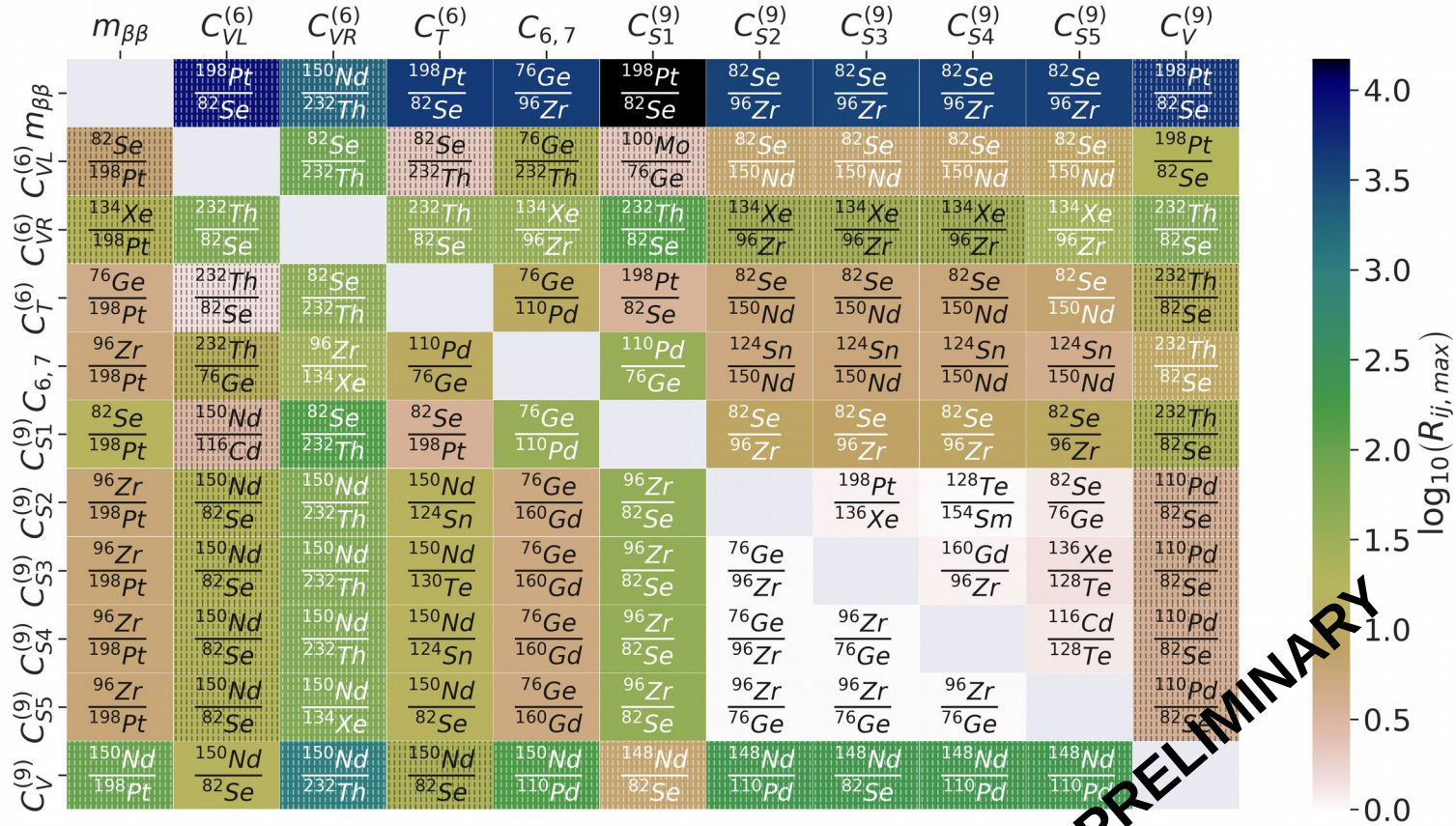
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Distinguishing $0\nu\beta\beta$ Mechanisms - Ratios



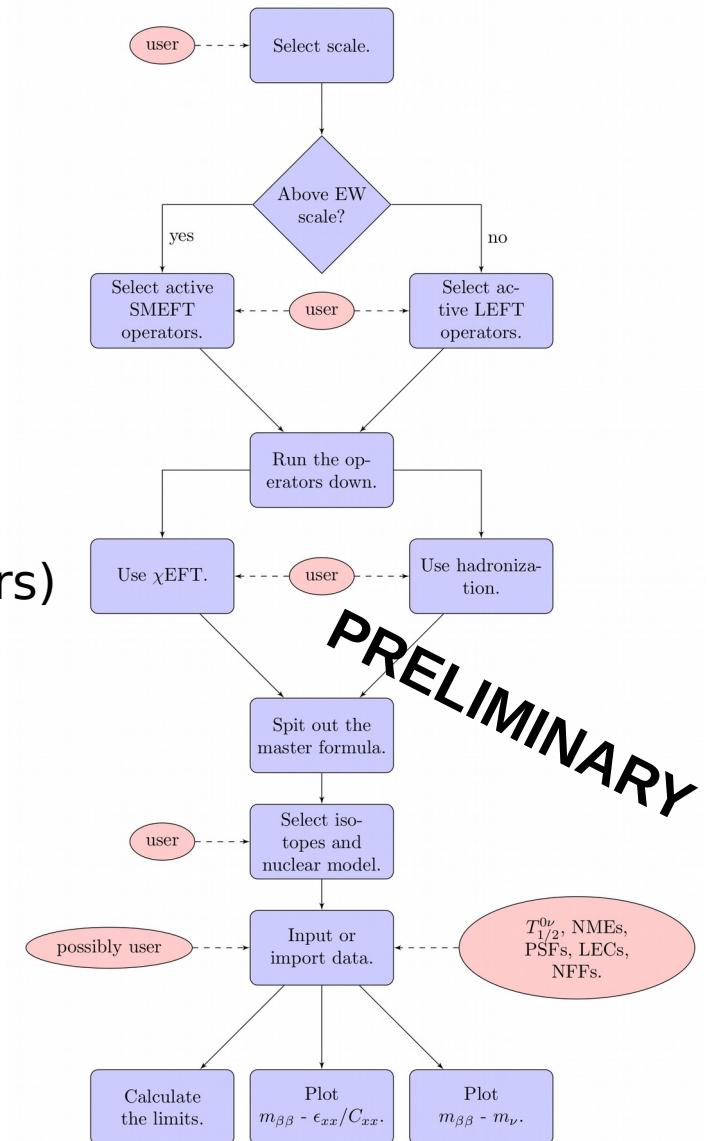
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$0\nu\beta\beta$ Automation Tool

- user inputs:
 - scale + selection of operators
 - isotope(s), type of NMEs
- data inputs:
 - nuclear matrix elements
 - phase-space factors
 - low-energy constants (and nuclear form factors)
- outputs:
 - half-life formula for the given case
 - limits on selected couplings
 - chosen contour plots showing correlations of different parameters
 - $m_{\beta\beta}$ vs. m_ν plots, etc.



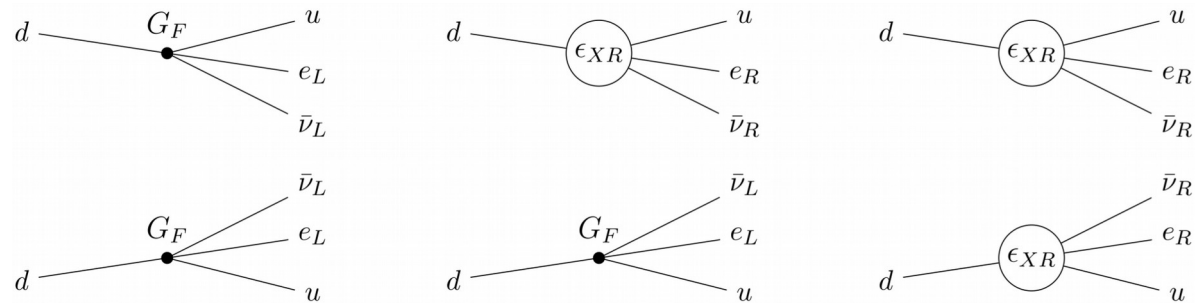
Exotic $2\nu\beta\beta$ – RHCs

- $2\nu\beta\beta$ decay in presence of right-handed currents?

$$\mathcal{L} = \frac{G_F \cos \theta_C}{\sqrt{2}} \left((1 + \delta_{\text{SM}} + \epsilon_{LL}) j_L^\mu J_{L\mu} + \epsilon_{RL} j_L^\mu J_{R\mu} + \epsilon_{LR} j_R^\mu J_{L\mu} + \epsilon_{RR} j_R^\mu J_{R\mu} \right) + \text{h.c.}$$

$$j_{L,R}^\mu = \bar{e} \gamma^\mu (1 \mp \gamma_5) \nu, \quad J_{L,R}^\mu = \bar{u} \gamma^\mu (1 \mp \gamma_5) d$$

- contributions:

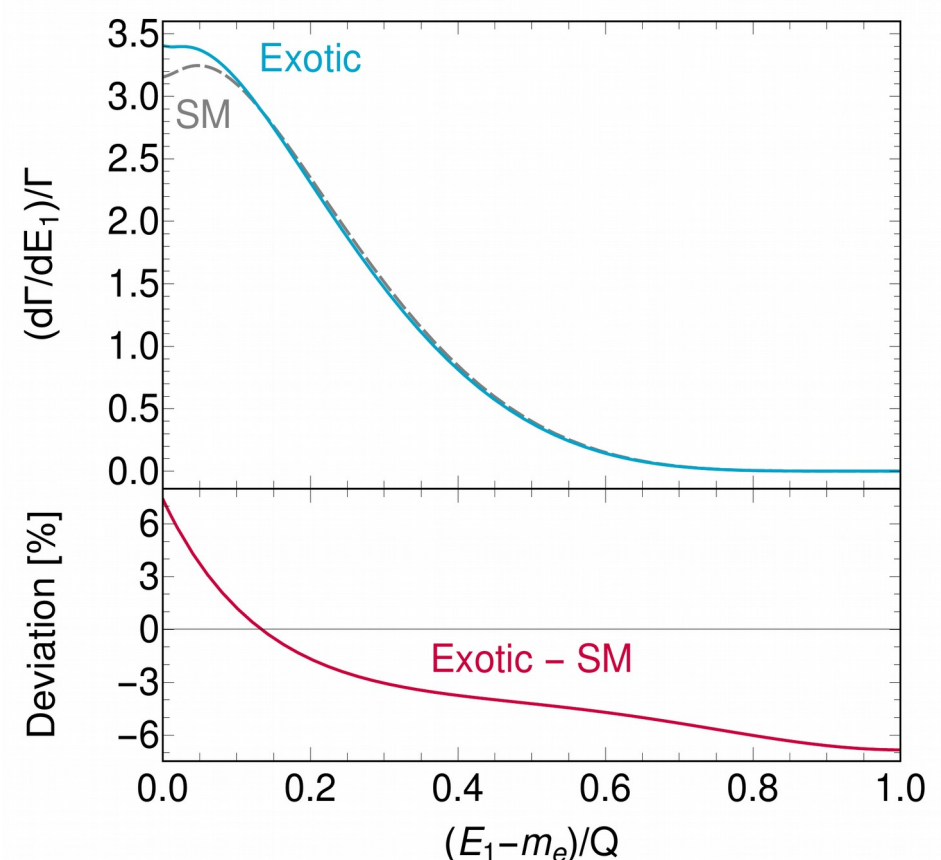
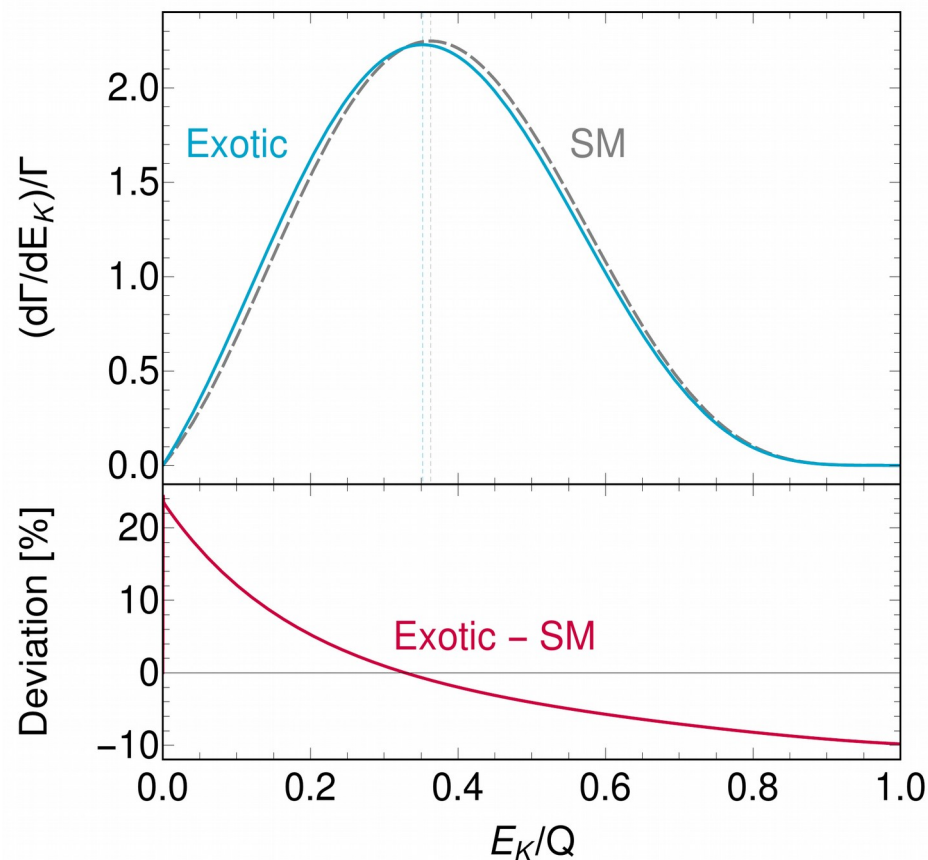


- take the one linear in exotic effective coupling ϵ_{XR} , calculate the observables and get the bound imposed by the collected experimental data

$$\rightarrow \text{rate: } [T_{1/2}^{2\nu\beta\beta}]^{-1} = \epsilon_{XR}^2 G_{2\nu\beta\beta} |M_{2\nu\beta\beta}|^2$$

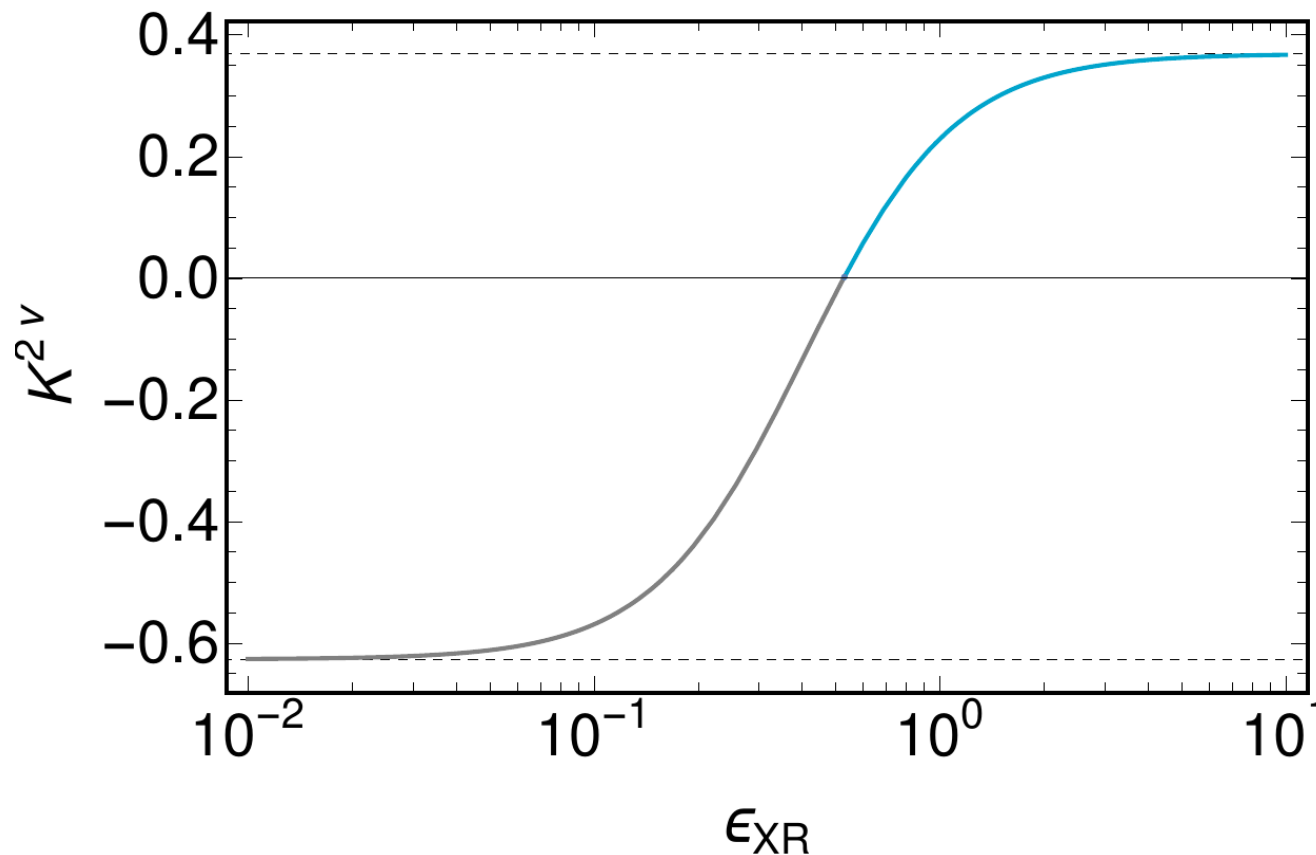
Electron Energy Distribution

- $2\nu\beta\beta$ decay (both the SM and the exotic RHC-induced contributions)
distribution in total kinetic (left) and single electron kinetic (right) energy



Electron Angular Correlation

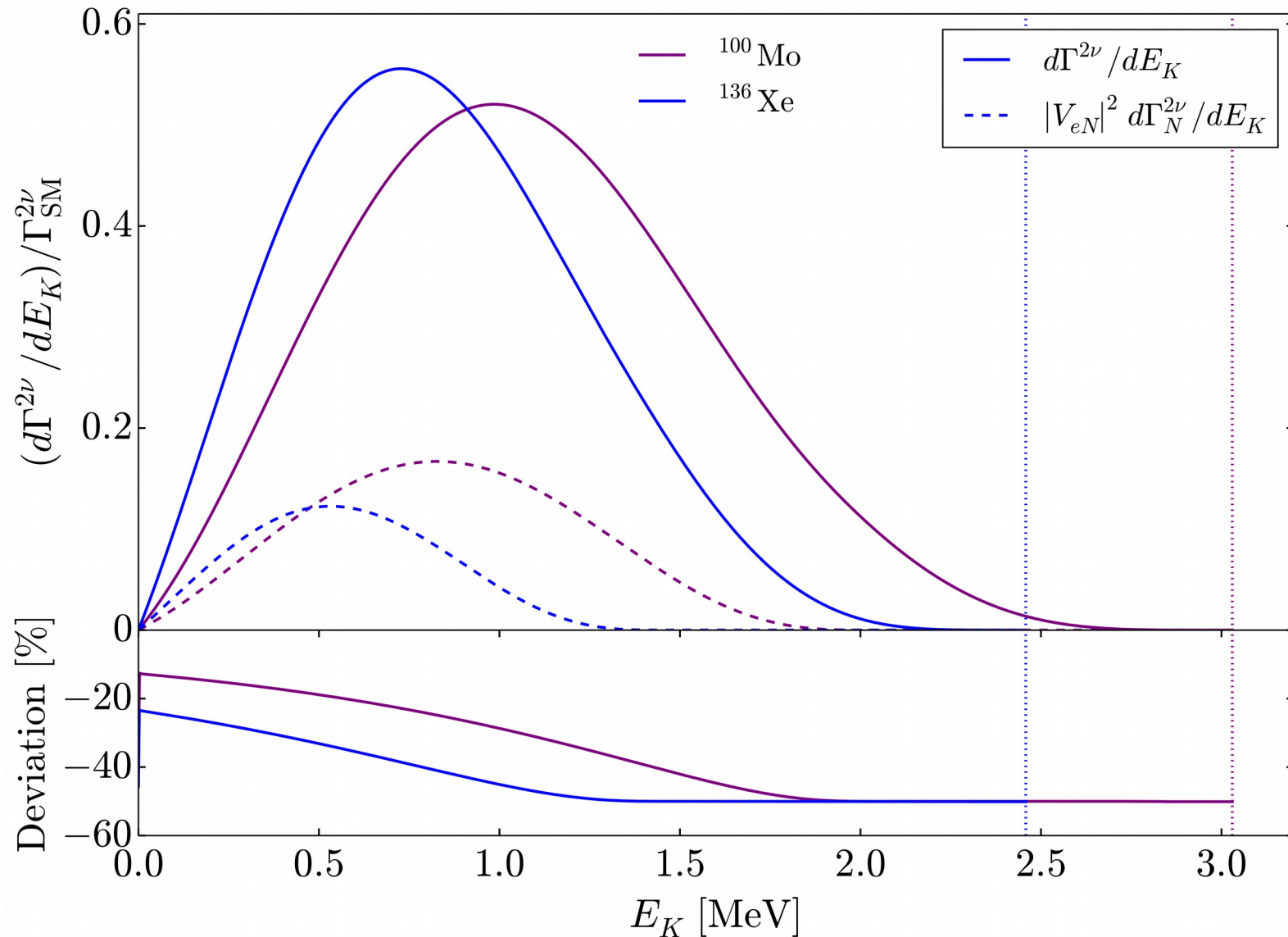
- observed angular correlation: mixture of the SM and exotic contributions – as a function of ϵ_{XR} interpolates between the SM ($\epsilon_{XR} = 0$) and exotic ($\epsilon_{XR} \gg 1$) cases



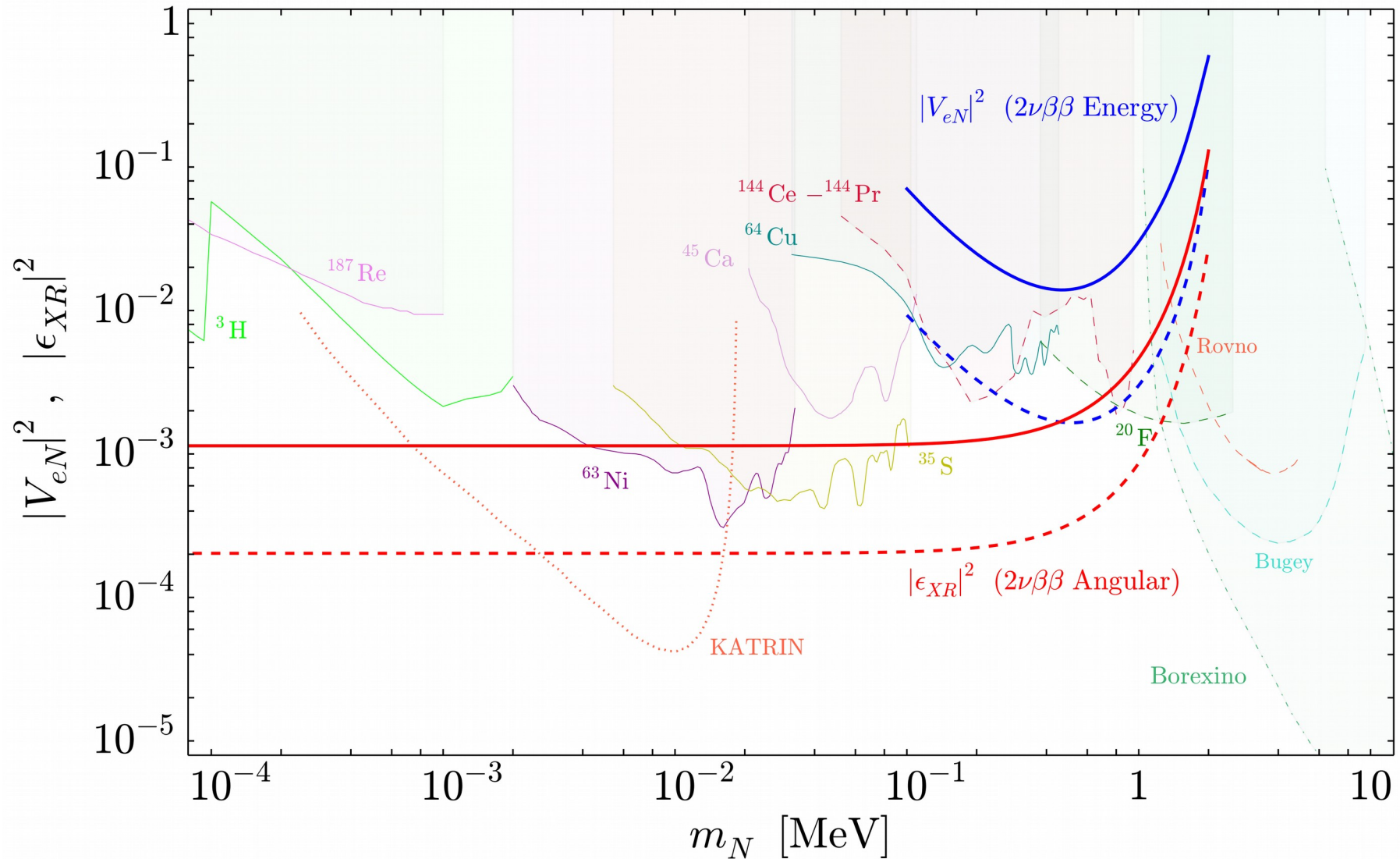
Bound on ϵ_{XR} Coupling

- angular distribution: $\frac{d\Gamma^{2\nu}}{d\cos\theta} = \frac{\Gamma^{2\nu}}{2} (1 + K^{2\nu} \cos\theta)$
with correlation factor: $K^{2\nu} \approx K_{SM}^{2\nu} + \alpha \epsilon_{XR}^2$
- forward-backward asymmetry: $A_{\theta}^{2\nu} = \frac{N_{\theta>\pi/2} - N_{\theta<\pi/2}}{N_{\theta>\pi/2} + N_{\theta<\pi/2}} = \frac{1}{2} K^{2\nu}$
- estimated accuracy of NEMO-3: $K_{SM}^{2\nu} = -0.63 \pm 0.0027$
→ bound on the effective coupling at 90% CL: $\epsilon_{XR} \lesssim 2.7 \times 10^{-2}$
- more stringent limit than the one obtained from the standard beta decay measurements
- SuperNEMO would further improve this bound
- rough estimates, dedicated experimental analysis necessary

Exotic $2\nu\beta\beta$ – Sterile Neutrino



Exotic $2\nu\beta\beta$ – Sterile Neutrino



Summary & Outlook

- A variety of different mechanisms may contribute to $0\nu\beta\beta$ and there are several possibilities how to distinguish among them; however, pinpointing the dominant contribution is not an easy task: measurements of energy spectra, angular correlation and usage of multiple isotopes necessary + LQCD input essential. The corresponding analysis to be published soon.
- Combining various contributions → involved, tedious calculations → $0\nu\beta\beta$ Automation Tool on the way → simple comparison of different mechanisms – mainly, calculation of corresponding limits and production of relevant plots. Anything else?
- Observation of $0\nu\beta\beta$ decay is the primary goal of double beta decay searches – no signal, yet, but we do see $2\nu\beta\beta$ decay, which can also probe New Physics!
→ Right-handed currents, sterile neutrinos, ν self-interactions...more?